# 2D/3D tidal modelling: T-UGOm spectral solver

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T-UGOm is basically a time-stepping model... so Why doing frequency-domain modeling ?

- First implemented to downscale global tidal atlases at open boundaries (currents issue)
- Realistic solutions
  - □ true turbulence closure
  - non-linearities solved by iterating solving cycle
- Extremely cheap compared to time-stepping modeling
- Easy implementation of various experiments/tuning allowed
- No separation issue (seasonal variability easier to examine)
- Well suited for frequency-domain data assimilation

Complementary to time-stepping simulation efforts...

## Spectral equations for astronomical tides

(equations for non-linear tides also available)

- Quasi-linearised, spectral Shallow-Water equations
  - Momentum equation
  - continuity equation

$$j\omega \mathbf{u} + \mathbf{f} \times \mathbf{u} = -g\nabla(\eta + \delta) + g\nabla\Pi - \mathbf{F}\mathbf{u} - \mathbf{D}\mathbf{u}$$
$$j\omega\eta + \nabla \cdot \mathbf{h}\mathbf{u} = \mathbf{0}$$

$$\mathbf{F} = \begin{bmatrix} \mathbf{r} & \mathbf{r}' \\ \mathbf{r}'' & \mathbf{r}''' \end{bmatrix} \qquad \mathbf{D} = c\rho_0 \frac{\kappa^{-1}}{\omega} \left[ \left( \mathcal{N}^2 - \omega^2 \right) \left( \omega^2 - \mathbf{f}^2 \right) \right]^{\frac{1}{2}} \left[ \nabla h \cdot^{\dagger} \nabla h \right]$$

see Lyard et al., 2006, Modelling the global ocean tides: modern insights from FES2004, Ocean Dynamics <u>http://dx.doi.org/10.1007/s10236-006-0086-x</u>

- Discrete equations (transport)
  - Momentum equation  $\mathbf{MU} = -\mathbf{G}\eta + \mathbf{F}$
  - Wave equation  $j\omega \mathbf{B}\eta \nabla \cdot \mathbf{M}^{-1}\mathbf{G}\eta = -\mathbf{D}\mathbf{M}^{-1}\mathbf{G}\eta = -\nabla \cdot \mathbf{M}^{-1}\mathbf{F}$

discrete velocity and elevation solution are fully consistent

- Double complex sparse matrix solver: PASTIX
  - OpenMP optimized, MPI available
  - Stable for large numbers of DoF (UMFPACK is not)

#### non-linearities: solved by iterating solving cycle



## Triangle elements 2D spectral modeling

#### Available discretisation





#### FES2014 discretisation, resolution, DoF

- □ 750 000 vertices
- □ 1 500 000 triangles
- □ 3 000 000 elevation nodes
- □ 4 500 000 velocity nodes









## 2D global ocean tides and storm surges modeling

#### **Objectives in preparation of SWOT mission**

- Continuously improve hydrodynamic tidal solution accuracy
- Use data assimilation (tide gauge and altimetryderived observations) to reach altimetry requirements
- Upgrade the global ocean storm surges simulator

wave	FES2012	FES2014	
M2	24 / 93 mm	13 / 53 mm	
S2	10 / 28 mm	8 / 18 mm	
K1	11 / 30 mm	9 / 23 mm	
O1	12 / 30 mm	7 / 20 mm	

Increase coastal resolution

comparisons against TP-J1-J2 (deep / shelf)



## **Spectral estuarine simulation**



Following damping and phase lags, Cd and H are modified in sections until fitting observations

So calibration works:

- Spectral calibration much quicker
- Still, automatized calibration schemes would be nice



Quadrangle elements 2D spectral modeling

#### **Pressure gradient issue**

- Pressure gradient #1 ill-constrains tangential gradient (ill-defined Corilolis term)
- Pressure gradient #2 provides a full gradient, but is over-lapping

Pressure gradient scheme #1 Identical to C-grid Pressure gradient scheme #2



#### Pressure gradient issue : impact on M2 tide simulation



Coriolis appropriate handling is essential...

#### HYCOM C-grid : Ushant Sea configuration



Required structured model inputs: T-grid, bathymetry, T-landmask

#### Some other examples



-0.02 0.02 0.06 0.1 0.14 0.18 0.22 0.26 0.3 0.34 0.38 0.42 0.46 0.5 0.54 0.58 0.62



- T-UGOm spectral solver works on multi-element, multidiscretisation, structured and unstructured grids
- Efficient to produce tidal solutions, from global to estuarine configuration
- Useful to downscale global ocean tidal atlases onto regional configuration (consitent currents at open boundaries)

## Triangle elements 3D spectral modeling

## **Vertical discretisation**

- Layer model
- Similar to generalized sigma-level discretisation
- Density is varying horizontally, but uniform along the vertical inside the 3D element



#### LGP0xLGP1, prismatic elements

- continuous, linear level displacements
- discontinuous, uniform velocities



## **Spectral 3D equations, basis**

- Navier-Stokes equations vertically integrated inside layers (similar to generalized sigma-level equations)
- Impermeable layers (no vertical fluxes)
- No horizontal diffusion, horizontal advection in sub-harmonic forcing
- Unknowns are 3D currents and level displacements



- 7th Coastal Altimetry Workshop – 7- 8 October 2013 -Boulder, CO -

## **Spectral 3D equations**

- Frequency-domain equations
- u: 3D velocity  $W = \varpi + \mathbf{V} \cdot \nabla_{H} \mathbf{S} + \frac{\partial \mathbf{S}}{\partial t} = \mathbf{V} \cdot \nabla_{H} \mathbf{S} + \frac{\partial \mathbf{S}}{\partial t}$ 1: level impermeability condition v: horizontal velocity (2D) V: horizontal transport w: vertical velocity  $i\omega\Delta\alpha + \nabla\cdot\overline{\Delta s}\mathbf{v} = i\omega\Delta\alpha + \nabla\cdot\mathbf{V} = \mathbf{0}$ 2: layer-integrated mass conservation +1 h: mean depth s: layer interface position (immersion)  $j\omega \mathbf{V} + \mathbf{f} \times \mathbf{V} = -\frac{1}{\rho} \int_{s_{1}}^{s_{1}} \nabla_{H} \mathbf{p} + \mathbf{g} \overline{\Delta s} \nabla (\Pi - \delta) + \left[ \kappa_{v} \frac{\partial \mathbf{v}}{\partial z} \right]_{s_{1}}^{s}$  $\alpha$ : layer interface (level) displacement 3: layer-integrated momentum equation  $\Delta s$ : laver thickness  $\Delta s$ : mean layer thickness  $\Delta \alpha$  : layer thickness change  $\kappa_{v} \frac{\partial \mathbf{V}}{\partial \mathbf{Z}}_{bottom} = -\mathbf{C}_{D} \| \mathbf{V} \| \mathbf{V}_{bottom}$ 4: bottom friction (vertical diffusion BC)  $\varpi$ : omega velocity  $\omega$ : tidal pulsation
  - Wave equation
    - Wave equation formed by substituting Eq. 3 in Eq. 2
    - Turbulence (and vertical diffusion) and non-linear terms solved by iterations

#### 3D barotropic turbulent scheme experiment





. Low resolution with high Kv can perform better than high resolution.

. The addition of a turbulent scheme (non constant Kv) can be highly beneficial in some cases.

#### 3D barotropic experiment



Vector difference (cm)	M2 - 2D	M2 - 3D (surface)	M2 - FES2012 (2D)
Tide gauges	3.6	3.4	3.6
Altimetry	2.1	2.0	2.5

#### Periodic channel test case: constant N=3.e-03, f=0



## **Spectral 3D baroclinic boundary conditions**

- Perfectly Matched Layers (PML):
  - Absorbing buffer zone (reflectionless under certain conditions)
  - Tricky to implement (barotropic forcing issue), in progress
- Sommerfeld radiative condition:
  - Effective for mono-chromatic waves
  - Choice : address 1<sup>st</sup> baroclinic mode
  - In consequence, vertical diffusivity is smoothly but strongly increased in a boundary buffer zone to minimize higher modes

$$\frac{\partial \eta'}{\partial t} + c \nabla_H \eta' \cdot \mathbf{n} = -j\omega\eta' + c \nabla_H \eta' \cdot \mathbf{n} = 0$$

radiative condition (at each level), variational solver

**n** : outward normal vector **s** : layer interface position (immersion) h : mean depth  $\eta$  : level (layer interface) displacement  $\eta_b = \frac{s+h}{h} \eta_{surface}$  : barotropic level displacement  $\eta' = \eta - \eta_b$  : baroclinic level displacement **c** : 1st baroclinic mode celerity  $\omega$  : tidal pulsation

### Instantaneous level displacement (t=0)



## (proper) Plane wave experiment



## (proper) Plane wave experiment



### Analytical solutions (2D and separated-modes 3D)

Tidal linearised, flat-bottom equations

$$\frac{\partial u}{\partial t} - fv = -g \frac{\partial \eta}{\partial x} - ru$$
$$\frac{\partial v}{\partial t} + fu = -g \frac{\partial \eta}{\partial y} - rv$$
$$\frac{\partial \eta}{\partial t} = -H\left(\frac{\partial u}{\partial x} + \frac{\partial v}{\partial x}\right)$$

Tidal frequency-domain ("spectral") equations

$$(-i\omega + r)\mathbf{u} + 2\vec{\Omega} \times \mathbf{u} = -g\nabla\eta$$
  

$$\frac{\partial\eta}{\partial t} + H\nabla \cdot \mathbf{u} = -i\omega\eta + H\nabla \cdot \mathbf{u} = 0$$
Dispersion relation:  

$$\mathbf{\Omega}' = \begin{bmatrix} 0\\ 0\\ \Omega\sin\varphi \end{bmatrix} \quad 2\mathbf{\Omega}' \times \mathbf{u} = 2\Omega \begin{bmatrix} -v\sin\varphi\\ u\sin\varphi\\ 0 \end{bmatrix} \quad \begin{pmatrix} (-i\omega + r)^2 + f^2 = \delta\\ c^2 = gH \end{bmatrix}$$
  

$$-\frac{c^2}{\delta} \Big[ (-i\omega + r)\nabla \cdot \nabla\eta + \frac{2f}{\delta}\nabla f \cdot ((-i\omega + r)\nabla\eta + 2\vec{\Omega}' \times \nabla\eta) \Big] = i\omega\eta$$

## Let's try pretty simple

- f taken as constant
- no friction

Dispersion relation reduces to:

$$\frac{\mathbf{c}^2}{\mathbf{f}^2 - \boldsymbol{\omega}^2} \nabla \cdot \nabla \eta = \eta$$

Plane waves: 
$$\eta(x, y, t) = \eta e^{i(k_x x + k_y y - \omega t)} \longrightarrow \nabla \cdot \nabla \eta = -(k_x^2 + k_y^2)\eta = -k_H^2 \eta$$

$$k_{H} = \frac{1}{c}\sqrt{\omega^{2} - f^{2}} \qquad \lambda = \frac{2\pi}{k_{H}} = \frac{\omega}{\sqrt{\omega^{2} - f^{2}}}cT$$

- Wavelength increase with latitude, inifinite at critical latitude
- > No solution above critical latitude
- > 1D wavelength at low latitudes  $k_H = \frac{\omega}{c}$   $\lambda = cT = \sqrt{gHT}$
- Not suitable to describe notable structures such as amphidromic points



## Let's try harder

- f taken as constant
- no friction

Dispersion relation reduces to:

$$\frac{\boldsymbol{c}^2}{\boldsymbol{f}^2 - \boldsymbol{\omega}^2} \nabla \cdot \nabla \eta = \eta$$

Kelvin waves: 
$$\eta(x, y, t) = \eta e^{-k_y y + i(k_x x - \omega t)}$$
  $\nabla \cdot \nabla \eta = -(k_x^2 - k_y^2)\eta$ 

$$k_{H} = k_{x} = \frac{1}{c}\sqrt{\omega^{2} - f^{2} + k_{y}^{2}} \qquad \lambda = \frac{2\pi}{k_{H}} = \frac{\omega}{\sqrt{\omega^{2} - f^{2} + k_{y}^{2}}} cT$$

- Coupling between across damping scale and wavelength
- > critical latitude is no more an issue
- > NOT 1D wavelength at low latitudes



## Let's try even harder: amphidromic point

- f taken as constant
- no friction

$$\frac{\boldsymbol{c}^2}{\boldsymbol{f}^2 - \boldsymbol{\omega}^2} \nabla \cdot \nabla \boldsymbol{\eta} = \boldsymbol{\eta}$$

Axial solution: 
$$\eta = \Phi(r)e^{i(\theta - \omega t)} \longrightarrow \frac{1}{\rho}\frac{\partial}{\partial\rho}\left(\rho\frac{\partial\eta}{\partial\rho}\right) - \frac{1}{\rho^2}\eta = -\frac{\left(\omega^2 - f^2\right)}{c^2}\eta = -\kappa^2\eta$$

It is a Sturm-Liouville problem, more exactly a Bessel differential equation

$$\eta = J_{\alpha=1}(\kappa\rho) = \sum_{m=0}^{+\infty} \frac{(-1)^m}{m! \, \Gamma(m+\alpha+1)} \left(\frac{x}{2}\right)^{2m+\alpha} = \sum_{m=0}^{+\infty} \frac{(-1)^m}{m! \, (m+2)!} \, e^{(2m+\alpha)\log\left(\frac{\kappa\rho}{2}\right)^{2m+\alpha}}$$

- Coupling between across scale and wavelength
- > critical latitude solution  $\eta = \eta_0 \frac{\rho}{\rho_0} e^{i(k_\theta \omega t)} = \eta_0 \frac{\rho}{\rho_0} e^{i(\theta \omega t)}$





- How much do we know about internal tides:
  - Temporal varibility
  - Horizontal scales (dispersion relation)
  - Is horizontal-vertical mode separation approach robust enough
  - Are higher modes predictable
  - Dissipation processes
- More analytical work needed (i.e. not only plane waves)
- How/where can we setup realistic configurations for modelling
  - Open boundary condition issues
  - Validation

## Summary

- Frequency-domain tidal dynamics solver is nested inside T-UGOm time-sequential solver
- Allows consistent downscaling for structured and unstructered models/configurations
- It solves for : 2D shallow-water, 2D + 1DV, 3D barotropic (shelf seas) and 3D baroclinic (internal tides)
- Accommodate different discretisation/element (generic coding)
- Frequency-domain data assimilation friendly
- Efficient for model parameters exploration
  - □ cheap
  - accurate
- On-going development (MPI) to address large domain 2D/3D simulations (10<sup>e+8</sup> dof)

#### Thank you for attention